

On a Class of Nonlinear Elliptic Systems Involving (p,q)-Laplacian

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Abstract. In this paper we study the existence of positive solutions for the system

$$\begin{cases} -\Delta_p u = \lambda f(u, v), & x \in \Omega, \\ -\Delta_q v = \lambda g(u, v), & x \in \Omega, \\ u = 0 = v, & x \in \partial\Omega, \end{cases}$$

where Δ_p is the so-called p-Laplacian operator i.e. $\Delta_p z = \operatorname{div}(|\nabla z|^{p-2} \nabla z)$, p and q are real numbers satisfying $1 < p, q < N$, λ is a real positive parameter, Ω is a bounded domain in R^N ($N \geq 1$) with smooth boundary $\partial\Omega$, and f, g are C^1 functions satisfying $\lim_{s \rightarrow \infty} f(s, t) = \infty = \lim_{t \rightarrow \infty} g(s, t)$, where each limit is uniform with respect to the other variable, $\lim_{|(s,t)| \rightarrow \infty} \frac{f(s,t)}{s^{p-1}} = \sigma$, and $\lim_{|(s,t)| \rightarrow \infty} \frac{g(s,t)}{t^{q-1}} = \delta$. In particular we do not assume any sign conditions on $f(0, 0)$ or $g(0, 0)$. For λ large we prove the existence of a large positive solution. Our approach is based on the method of sub-super.

Key Words and Phrases: Elliptic system; Positive solutions

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1. Introduction

The aim of this article is to study the existence of positive solutions for some nonlinear elliptic systems of the form

$$\begin{cases} -\Delta_p u = \lambda f(u, v), & x \in \Omega, \\ -\Delta_q v = \lambda g(u, v), & x \in \Omega, \\ u = 0 = v, & x \in \partial\Omega, \end{cases} \quad (1)$$

here Δ_p is the so-called p-Laplacian operator i.e. $\Delta_p z = \operatorname{div}(|\nabla z|^{p-2} \nabla z)$, p and q are real numbers satisfying $1 < p, q < N$, λ is a real positive parameter, Ω is a bounded domain in R^N ($N \geq 1$) with smooth boundary $\partial\Omega$, and $f, g : [0, \infty) \times [0, \infty) \rightarrow R$ are C^1 functions satisfying the following assumptions:

(A1) $\lim_{s \rightarrow \infty} f(s, t) = \infty = \lim_{t \rightarrow \infty} g(s, t)$, where each limit is uniform with respect to the other variable;

(A2) $\lim_{|(s,t)| \rightarrow \infty} \frac{f(s,t)}{s^{p-1}} = \sigma$, and $\lim_{|(s,t)| \rightarrow \infty} \frac{g(s,t)}{t^{q-1}} = \delta$;

(A3) $s \mapsto f(s, t)$ and $s \mapsto g(s, t)$ are nondecreasing for every $t > 0$;

(A4) $t \mapsto f(s, t)$ and $t \mapsto g(s, t)$ are nondecreasing for every $s > 0$.

Let $\zeta_1(x)$, $\zeta_2(x)$ be the positive solutions, respectively, of the problems

$$\begin{cases} -\Delta_p \zeta_1 = 1, & x \in \Omega, \\ \zeta_1 = 0, & x \in \partial\Omega, \end{cases}$$

and

$$\begin{cases} -\Delta_q \zeta_2 = 1, & x \in \Omega, \\ \zeta_2 = 0, & x \in \partial\Omega, \end{cases}$$

where Ω is as before. Let $l_1 = \|\zeta_1\|_\infty$, $l_2 = \|\zeta_2\|_\infty$, and we assume that

$$\sigma < \frac{1}{l_1^{p-1}}, \quad \delta < \frac{1}{l_2^{q-1}}, \quad (2)$$

where σ and δ is in (A2).

The problem (1) arises in the theory of quasiregular and quasiconformal mappings or in the study of non-Newtonian fluids [2]. In the later case the quantity p is a characteristic of the medium. Media with $p > 2$ are called dilatant fluid and those with $p < 2$ are called pseudoplastics.

The solvability of system (1) has been studied by various methods, fibering [4], bifurcation [9], via the mountain pass theorem [3]. See [1, 5] where the authors discussed the system (1) when $p = q = 2$, $f(u, v) = \tilde{f}(u)$, $g(u, v) = \tilde{g}(u)$, \tilde{f}, \tilde{g} are increasing and $\tilde{f}, \tilde{g} \geq 0$. In [11], the authors extended the study of [5], to the case when no sign conditions on $f(0)$ or $g(0)$ were required and in [12] they extend this study to the case when $p = q > 1$. Here we focus on further extending the study in [12] for the quasilinear elliptic systems with much stronger coupling. Due to this strong coupling conditions, the extensions are challenging and nontrivial. Our approach is based on the method of sub-super solutions, see [8]. We refer to [6, 7, 10, 13] for additional results on elliptic problems.

2. Existence results

To prove our existence results we use the method of sub-super solutions. To do so, we now define sub and super solutions of (1). Let $W_0^{1,s} = W_0^{1,s}(\Omega)$, $s > 1$, denote the usual Sobolev space.

Definition 1. A pair of nonnegative functions (ψ_1, ψ_2) , (z_1, z_2) in $W_0^{1,p} \times W_0^{1,q}$ are called a weak subsolution and supersolution of (1) if they satisfy $\psi_i(x) \leq z_i(x)$ in Ω for $i = 1, 2$, and

$$\int_{\Omega} |\nabla \psi_1|^{p-2} \nabla \psi_1 \nabla w_1 dx \leq \lambda \int_{\Omega} f(\psi_1, \psi_2) w_1 dx, \quad (3)$$

$$\int_{\Omega} |\nabla \psi_2|^{q-2} \nabla \psi_2 \nabla w_2 dx \leq \lambda \int_{\Omega} g(\psi_1, \psi_1) w_2 dx, \quad (4)$$

$$\int_{\Omega} |\nabla z_1|^{p-2} \nabla z_1 \nabla w_1 dx \geq \lambda \int_{\Omega} f(z_1, z_2) w_1 dx, \quad (5)$$

and

$$\int_{\Omega} |\nabla z_2|^{q-2} \nabla z_2 \nabla w_2 dx \geq \lambda \int_{\Omega} g(z_1, z_2) w_2 dx, \quad (6)$$

for all $w_1(x) \in W_0^{1,p}$, $w_2(x) \in W_0^{1,q}$, with $w_1, w_2, \geq 0$.

We shall obtain the existence of positive solution to system (1) by constructing a positive subsolution (ψ_1, ψ_2) and supersolution (z_1, z_2) .

Our main result is formulate in the following theorem.

Theorem 1. Assume that hypotheses (A1)–(A4) and (2) hold. Then there exists a positive number λ_0 such that (1) has a large positive solution (u, v) for $\lambda > \lambda_0$.

Proof. Let λ_1 and λ_2 be the first eigenvalue of the problems, respectively,

$$\begin{aligned} -\Delta_p \phi_1 &= \lambda_1 \phi_1^{p-1}, & x \in \Omega, & \quad \phi_1 = 0, & \quad x \in \partial\Omega, \\ -\Delta_q \phi_2 &= \lambda_2 \phi_2^{q-1}, & x \in \Omega, & \quad \phi_2 = 0, & \quad x \in \partial\Omega, \end{aligned}$$

where ϕ_1 and ϕ_2 denote the corresponding positive eigenfunctions, respectively, satisfying $\|\phi_i\|_{\infty} = 1$ for $i = 1, 2$. By (A1) we can take $a_1, a_1 > 0$ such that

$$f(s, t) > -a_1, \quad g(s, t) > -a_2,$$

for all $s, t \geq 0$. Let $b_1, b_2 > 0$ be such that

$$\lambda_1 \phi_1^p - |\nabla \phi_1|^p \leq -b_1, \quad \lambda_2 \phi_2^q - |\nabla \phi_2|^q \leq -b_2, \quad x \in \bar{\Omega}_\eta,$$

where $\bar{\Omega}_\eta = \{x \in \Omega \mid \text{dist}(x, \partial\Omega) \leq \eta\}$. (This is possible since $|\nabla \phi_i|^r \neq 0$ on $\partial\Omega$ while $\phi_i = 0$ on $\partial\Omega$ for $r = p, q$ and $i = 1, 2$). We shall verify that

$$(\psi_1, \psi_2) = \left(\left(\frac{\lambda a_1}{b_1} \right)^{\frac{1}{p-1}} \left(\frac{p-1}{p} \right) \phi_1^{\frac{p}{p-1}}, \left(\frac{\lambda a_2}{b_2} \right)^{\frac{1}{q-1}} \left(\frac{q-1}{q} \right) \phi_2^{\frac{q}{q-1}} \right),$$

is a subsolution of (1) for λ large. Let the test function $w_1(x) \in W_0^{1,p}$, with $w_1 \geq 0$. Then it follows from (3) that

$$\begin{aligned} \int_{\Omega} |\nabla \psi_1|^{p-2} \nabla \psi_1 \cdot \nabla w_1 &= \frac{\lambda a_1}{b_1} \int_{\Omega} \phi_1 |\nabla \phi_1|^{p-2} \nabla \phi_1 \nabla w_1 dx \\ &= \frac{\lambda a_1}{b_1} \left\{ \int_{\Omega} |\nabla \phi_1|^{p-2} \nabla \phi_1 \nabla (\phi_1 w_1) dx - \int_{\Omega} |\nabla \phi_1|^p w_1 dx \right\} \\ &= \frac{\lambda a_1}{b_1} \int_{\Omega} (\lambda_1 \phi_1^p - |\nabla \phi_1|^p) w_1 dx. \end{aligned}$$

Since on $\bar{\Omega}_\eta$ we have $\lambda_1 \phi_1^p - |\nabla \phi_1|^p \leq -b_1$, which implies that

$$\frac{a_1}{b_1} (\lambda_1 \phi_1^p - |\nabla \phi_1|^p) \leq f(\psi_1, \psi_2).$$

It is well known that $\frac{\partial \phi_i}{\partial n} < 0$ on $\partial\Omega$ where n is the unit outward normal for $i = 1, 2$. Hence there exists $\beta > 0$ such that $\phi_1 \geq \beta > 0$ in $\Omega_0 = \Omega \setminus \bar{\Omega}_\eta$. Therefore, from (A1) for λ large, we have

$$\frac{a_1}{b_1} (\lambda_1 \phi_1^p - |\nabla \phi_1|^p) \leq \frac{a_1}{b_1} \lambda_1 \leq f(\psi_1, \psi_2).$$

Hence

$$\int_{\Omega} |\nabla \psi_1|^{p-2} \nabla \psi_1 \nabla w_1 dx \leq \lambda \int_{\Omega} f(\psi_1, \psi_2) w_1 dx.$$

Similarly, we have

$$\int_{\Omega} |\nabla \psi_2|^{q-2} \nabla \psi_2 \nabla w_2 dx \leq \lambda \int_{\Omega} g(\psi_1, \psi_1) w_2 dx,$$

for all $w_2(x) \in W_0^{1,q}$, with $w_2 \geq 0$. Thus, (ψ, ψ) is a subsolution of (1).

Next, we construct a supersolution (z_1, z_2) of (1). We denote

$$(z_1, z_2) = (\lambda^{\frac{1}{p-1}} C_1 \zeta_1, \lambda^{\frac{1}{q-1}} C_2 \zeta_2),$$

where ζ_1, ζ_2 are as before, and $C_1, C_2 > 0$ are large enough such that

$$\lambda f(\lambda^{\frac{1}{p-1}} C_1 l_1, \lambda^{\frac{1}{q-1}} C_2 l_2) \leq (C_1 \lambda^{1/p-1})^{p-1}, \quad (7)$$

and

$$\lambda g(\lambda^{\frac{1}{p-1}} C_1 l_1, \lambda^{\frac{1}{q-1}} C_2 l_2) \leq (C_2 \lambda^{1/q-1})^{q-1}. \quad (8)$$

Here (7) and (8) are possible by (A2), and (2). We shall verify that (z_1, z_2) is a supersolution of (1). To this end, let $w_1(x) \in W_0^{1,p}$, $w_2(x) \in W_0^{1,q}$, with $w_1, w_2 \geq 0$. Then we obtain from (7), (8), (A3) and (A4), that

$$\begin{aligned} \int_{\Omega} |\nabla z_1|^{p-2} \nabla z_1 \nabla w_1 dx &= \lambda C_1^{p-1} \int_{\Omega} |\nabla \zeta_1|^{p-2} \nabla \zeta_1 \nabla w_1 dx \\ &= \lambda C_1^{p-1} \int_{\Omega} w_1 dx \\ &\geq \lambda \int_{\Omega} f(\lambda^{\frac{1}{p-1}} C_1 l_1, \lambda^{\frac{1}{q-1}} C_2 l_2) w_1 dx \\ &\geq \lambda \int_{\Omega} f(\lambda^{\frac{1}{p-1}} C_1 \zeta_1, \lambda^{\frac{1}{q-1}} C_2 \zeta_2) w_1 dx \\ &= \lambda \int_{\Omega} f(z_1, z_2) w_1 dx. \end{aligned}$$

Similarly we have

$$\int_{\Omega} |\nabla z_2|^{q-2} \nabla z_2 \nabla w_2 dx \geq \lambda \int_{\Omega} g(z_1, z_2) w_2 dx.$$

Thus, (z_1, z_2) is a supersolution of (1) with $z_i \geq \psi_i$ in Ω for large C_1, C_2 , $i = 1, 2$. Thus, by the comparison principle, there exists a solution (u, v) of (1) with $\psi_1 \leq u \leq z_1$, $\psi_2 \leq v \leq z_2$. This completes the proof of Theorem 1. ◀

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